

L^1 averaging lemma for transport equations with Lipschitz force fields

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Abstract

The purpose of this note is to extend the L^1 averaging lemma of Golse and Saint-Raymond [10] to the case of a kinetic transport equation with a force field $F(x) \in W^{1,\infty}$. To this end, we will prove a local in time mixing property for the transport equation $\partial_t f + v \cdot \nabla_x f + F \cdot \nabla_v f = 0$.

Introduction

Let $d \in \mathbb{N}^*$ and $1 < p < +\infty$. We consider \mathbb{R}^d equipped with the Lebesgue measure. Let $f(x, v)$ and $g(x, v)$ be two measurable functions in $L^p(\mathbb{R}^d \times \mathbb{R}^d)$ satisfying the transport equation:

$$v \cdot \nabla_x f = g. \quad (0.1)$$

Although transport equations are of hyperbolic nature (and thus there is a priori no regularizing effect), it was first observed for by Golse, Perthame and Sentis in [8] and then by Golse, Lions, Perthame and Sentis [9] (see also Agoshkov [1] for related results obtained independently) that the velocity average (or moment) $\rho(x) = \int f \Psi(v) dv$ with $\Psi \in \mathcal{C}_c^\infty(\mathbb{R}^d)$ is smoother than f and g : more specifically it belongs to some Sobolev space $W^{s,p}(\mathbb{R}^d)$ with $s > 0$. These kinds of results are referred to as "velocity averaging lemma". The analogous results in the time-dependent setting also hold, that is for the equation:

$$\partial_t f + v \cdot \nabla_x f = g. \quad (0.2)$$

Refined results with various generalizations (like derivatives in the right-hand side, functions with different integrability in x and v ...) were obtained in [7], [3], [14], [13]. There exist many other interesting contributions. We refer to Jabin [12] which is a rather complete review on the topic.

Velocity averaging lemmas are tools of tremendous importance in kinetic theory since they provide some strong compactness which is very often necessary to study non-linear terms (for instance when one considers an approximation scheme to build weak solutions, or for the study of asymptotic regimes). There are numerous applications of these lemmas; two emblematic results are the existence of renormalized solutions to the Boltzmann equation [5] and the existence of global weak solutions to the Vlasov-Maxwell system [6]. Both are due to DiPerna and Lions.

The limit case $p = 1$ is actually of great interest. In general, for a sequence (f_n) uniformly bounded in $L^1(dx \otimes dv)$ with $v \cdot \nabla_x f_n$ also uniformly bounded in $L^1(dx \otimes dv)$, the sequence of velocity averages $\rho_n = \int f_n \Psi(v) dv$ is not relatively compact in $L^1(dx)$

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(we refer to [9] for an explicit counter-example). This lack of compactness is due to the weak compactness pathologies of L^1 . Indeed, as soon as we add some weak compactness to the sequence (or equivalently some equiintegrability in x and v in view of the classical Dunford-Pettis theorem), then we recover some strong compactness in L^1 for the moments (see Proposition 3 of [9] or Proposition 3 below).

We recall precisely the notion of equiintegrability which is central in this paper.

Definition. 1. (*Local equiintegrability in x and v*)

Let (f_ϵ) be a bounded family of $L^1_{loc}(dx \otimes dv)$. It is said locally equiintegrable in x and v if and only if for any $\eta > 0$ and for any compact subset $K \subset \mathbb{R}^d \times \mathbb{R}^d$, there exists $\alpha > 0$ such that for any measurable set $A \subset \mathbb{R}^d \times \mathbb{R}^d$ with $|A| < \alpha$, we have for any ϵ :

$$\int_A \mathbb{1}_K(x, v) |f_\epsilon(x, v)| dv dx \leq \eta. \quad (0.3)$$

2. (*Local equiintegrability in v*)

Let (f_ϵ) be a bounded family of $L^1_{loc}(dx \otimes dv)$. It is said locally equiintegrable in v if and only if for any $\eta > 0$ and for any compact subset $K \subset \mathbb{R}^d \times \mathbb{R}^d$, there exists $\alpha > 0$ such that for each family $(A_x)_{x \in \mathbb{R}^d}$ of measurable sets of \mathbb{R}^d satisfying $\sup_{x \in \mathbb{R}^d} |A_x| < \alpha$, we have for any ϵ :

$$\int \left(\int_{A_x} \mathbb{1}_K(x, v) |f_\epsilon(x, v)| dv \right) dx \leq \eta. \quad (0.4)$$

We observe that local equiintegrability in (x, v) always implies local equiintegrability in v , whereas the converse is false in general.

The major improvement of the paper of Golse and Saint-Raymond [10] is to show that actually, only equiintegrability in v is needed to obtain the L^1 compactness for the moments. This observation was one of the key arguments of their outstanding paper [11] which establishes the convergence of renormalized solutions to the Boltzmann equation in the sense of DiPerna-Lions to weak solutions to the Navier-Stokes equation in the sense of Leray.

More precisely, the result they prove is Theorem 1 stated afterwards, with $F = 0$ (free transport case). The aim of this paper is to show that the result also holds if one adds some force field $F(x) = (F_i(x))_{1 \leq i \leq d}$ with $F \in W^{1, \infty}(\mathbb{R}^d)$:

Theorem 1. *Let (f_ϵ) be a family bounded in $L^1_{loc}(dx \otimes dv)$ locally equiintegrable in v and such that $v \cdot \nabla_x f_\epsilon + F \cdot \nabla_v f_\epsilon$ is bounded in $L^1_{loc}(dx \otimes dv)$. Then :*

1. (f_ϵ) is locally equiintegrable in both variables x **and** v .
2. For all $\Psi \in \mathcal{C}^1_c(\mathbb{R}^d)$, the family $\rho_\epsilon(x) = \int f_\epsilon(x, v) \Psi(v) dv$ is relatively compact in $L^1_{loc}(dx)$.

One key ingredient of the proof for $F = 0$ is the nice dispersion properties of the free transport operator. We will show in Section 2 that an analogue also holds for small times when $F \neq 0$:

Proposition 1. *Let $F(x)$ be a Lipschitz vector field. There exists a maximal time $\tau > 0$ (depending only on $\|\nabla_x F\|_{L^\infty}$) such that, if f is the solution to the transport equation:*

$$\begin{cases} \partial_t f + v \cdot \nabla_x f + F \cdot \nabla_v f = 0, \\ f(0, \cdot, \cdot) = f^0 \in L^p(dx \otimes dv), \end{cases}$$

Then:

$$\forall |t| \leq \tau, \|f(t)\|_{L_x^\infty(L_v^1)} \leq \frac{2}{|t|^d} \|f^0\|_{L_x^1(L_v^\infty)}. \quad (0.5)$$

Let us also mention that the main theorem generalizes to the time-dependent setting, for transport equations of the form (0.2). The usual trick to deduce such a result from the stationary case is to enlarge the phase space. Indeed we can consider $x' = (t, x)$ in \mathbb{R}^{d+1} endowed with the Lebesgue measure, and $v' = (t, v)$ in \mathbb{R}^{d+1} endowed with the measure $\mu = \delta_{t=1} \otimes \text{Leb}$ (where δ is the dirac measure). Then such a measure μ satisfies property (2.1) of [9]. As a consequence, all the results of Section 1 will still hold.

Nevertheless, we observe that our key local in time mixing estimate (0.5) seems to not hold when \mathbb{R}^{d+1} is equipped with the new measure μ (the main problem being that the only speed associated to the first component of v' is 1). For this reason, we can not prove that equiintegrability in v implies equiintegrability in t . One result (among other possible variants) is the following:

Theorem 2. *Let $F(t, x) \in \mathcal{C}^0(\mathbb{R}^+, W^{1,\infty}(\mathbb{R}^d))$. Let (f_ϵ) be a family bounded in $L_{loc}^1(dt \otimes dx \otimes dv)$ locally equiintegrable in v and such that $\partial_t f_\epsilon + v \cdot \nabla_x f_\epsilon + F \cdot \nabla_v f_\epsilon$ is bounded in $L_{loc}^1(dt \otimes dx \otimes dv)$. Then :*

1. (f_ϵ) is locally equiintegrable in the variables x and v (but not necessarily with respect to t).
2. For all $\Psi \in \mathcal{C}_c^1(\mathbb{R}^d)$, the family $\rho_\epsilon(t, x) = \int f_\epsilon(t, x, v) \Psi(v) dv$ is relatively compact with respect to the x variable in $L_{loc}^1(dt \otimes dx)$, that is, for any compact $K \subset \mathbb{R}_t^+ \times \mathbb{R}_x^d$:

$$\lim_{\delta \rightarrow 0} \sup_{\epsilon} \sup_{|x'| \leq \delta} \|(\mathbb{1}_K \rho_\epsilon)(t, x + x') - (\mathbb{1}_K \rho_\epsilon)(t, x)\|_{L^1(dt \otimes dx)} = 0. \quad (0.6)$$

Another possibility is to assume that f_ϵ is bounded in $L_{t,loc}^\infty(L_{x,v,loc}^1)$, in which case we will get equiintegrability in t, x and v and thus compactness for ρ_ϵ in t and x . We refer to [11], Lemma 3.6, for such a statement in the free transport case.

Remarks 0.1. 1. *Since the result of Theorem 1 is essentially of local nature, we could slightly weaken the assumption on F :*

For any $R > 0$,

$$\exists M(R), \forall |x_1|, |x_2| \leq R, \quad |F(x_1) - F(x_2)| \leq M(R)|x_1 - x_2|. \quad (0.7)$$

In other words we can deal with $F \in W_{loc}^{1,\infty}$.

2. *With the same proof, we can treat the case of force fields $F(x, v) \in W_{x,v,loc}^{1,\infty}$ with zero divergence in v :*

$$\text{div}_v F = 0.$$

Typically we may think of the Lorentz force $v \wedge B$ where B is a smooth magnetic field.

3. *We can handle a family of force fields (F_ϵ) depending on ϵ as soon as (F_ϵ) is uniformly bounded in $W^{1,\infty}(\mathbb{R}^d)$.*

The following of the paper is devoted to the proof of Theorem 1. In Section 1, we prove that a family satisfying the assumptions of Theorem 1 and in addition locally equiintegrable in x and v , has moments which are relatively strongly compact in L^1 . In Section 2, we investigate the local in time mixing properties of the transport equation $\partial_t f + v \cdot \nabla_x f + F \cdot \nabla_v f = 0$. Finally in the last section, thanks to the mixing properties we establish, we show by an interpolation argument that equiintegrability in v provides some equiintegrability in x .

1 A first step towards L^1 compactness

The first step is to show that under the assumptions of Theorem 1, point 1 implies point 2. Using classical averaging lemma in L^2 ([6], [7]), we first prove the following L^2 averaging lemma.

Lemma 1. *Let $f, g \in L^2(dx \otimes dv)$ satisfy the transport equation:*

$$v \cdot \nabla_x f + F \cdot \nabla_v f = g. \quad (1.1)$$

Then for all $\Psi \in C_c^1(\mathbb{R}^d)$, $\rho(x) = \int f(x, v) \Psi(v) dv \in H_x^{1/4}$. Moreover,

$$\|\rho\|_{H_x^{1/4}} \leq C \left(\|F\|_{L_x^\infty} \|f\|_{L_{x,v}^2} + \|g\|_{L_{x,v}^2} \right). \quad (1.2)$$

(C is a constant depending only on Ψ .)

Proof. The standard idea is to consider $-F \cdot \nabla_v f + g$ as a source. Then, since $\operatorname{div}_v F(x) = 0$, we have :

$$-F \cdot \nabla_v f + g = - \sum_{i=1}^d \frac{\partial}{\partial v_i} (F_i f) + g.$$

We conclude by applying the L^2 averaging lemma of [6], Theorem 3. \square

We recall now in Proposition 2 an elementary and classical representation result, obtained by the method of characteristics.

Let $b = (v, F)$, $Z = (X, V)$. Since $F \in W^{1,\infty}$, b satisfies the hypotheses of the global Cauchy-Lipschitz theorem. We therefore consider the trajectories defined by:

$$\begin{cases} Z'(t; x_0, v_0) = b(Z(t; x_0, v_0)) \\ Z(0; x_0, v_0) = (x_0, v_0). \end{cases} \quad (1.3)$$

For all time, the application $(x_0, v_0) \mapsto Z(t; x_0, v_0) = (X(t; x_0, v_0), V(t; x_0, v_0))$ is well-defined and is a C^1 diffeomorphism. Moreover, since b does not depend explicitly on time, it is also classical that $Z(t)$ is a group. The inverse is thus given by $(x, v) \mapsto Z(-t; x, v)$.

Remark 1. *Since $\operatorname{div}(b) = 0$, Liouville's theorem shows that the volumes in the phase space are preserved (the jacobian determinant of Z is equal to 1).*

Proposition 2. *1. The time-dependent Cauchy problem :*

$$\begin{cases} \partial_t f + v \cdot \nabla_x f + F \cdot \nabla_v f = 0, \\ f(0, \cdot, \cdot) = f^0 \in L^p(dx \otimes dv) \end{cases} \quad (1.4)$$

has a unique solution (in the distributional sense) represented by

$$f(t, x, v) = f^0(X(-t; x, v), V(-t; x, v)) \in L^p(dx \otimes dv).$$

2. For any $\lambda > 0$, the transport equation

$$\lambda f(x, v) + v \cdot \nabla_x f + F \cdot \nabla_v f = g \in L^p(dx \otimes dv) \quad (1.5)$$

has a unique solution (in the distributional sense) represented by:

$$R_\lambda : g(x, v) \mapsto f(x, v) = \int_0^{+\infty} e^{-\lambda s} g(X(-s; x, v), V(-s; x, v)) ds \in L^p(dx \otimes dv).$$

In addition, R_λ is a linear continuous map on L^p with a norm equal to $\frac{1}{\lambda}$.

Using Rellich's compactness theorem, we straightforwardly have the following corollary:

Corollary 1. *The linear continuous map $T_{\lambda, \Psi} :$*

$$L^2(dx \otimes dv) \rightarrow L^2_{loc}(dx)$$

$$g \mapsto \rho = \int R_\lambda(g)(\cdot, v) \Psi(v) dv$$

is compact for all $\Psi \in \mathcal{C}_c^1(\mathbb{R}^d)$ and all $\lambda > 0$.

Proof. Using Lemma 1 and Proposition 2, we have:

$$\|T_{\lambda, \Psi}(g)\|_{H_x^{1/4}} \leq C(1 + \|F\|_{L^\infty}) \|g\|_{L^2_{x,v}}.$$

The conclusion follows. \square

Using this compactness property, as in Proposition 3 of [9], we can show the next result:

Proposition 3. *Let \mathcal{K} be a bounded subset of $L^1(dx \otimes dv)$ equiintegrable in x and v (in view of the Dunford-Pettis theorem, it means in other words that \mathcal{K} is weakly compact in L^1), then $T_{\lambda, \Psi}(\mathcal{K})$ is relatively strongly compact in $L^1_{loc}(dx)$.*

Proof. We recall the proof of this result for the sake of completeness.

The proof is based on a real interpolation argument. We fix a parameter $\eta > 0$. For any $g \in \mathcal{K}$ and any $\alpha > 0$, we may write :

$$g = g_1^\alpha + g_2^\alpha,$$

with

$$\begin{aligned} g_1^\alpha &= \mathbb{1}_{\{|g(x,v)| > \alpha\}} g, \\ g_2^\alpha &= \mathbb{1}_{\{|g(x,v)| \leq \alpha\}} g. \end{aligned}$$

Then, by linearity of $T_{\lambda, \Psi}$, we write $u = T_{\lambda, \Psi}(g) = u_1 + u_2$, with $u_1 = T_{\lambda, \Psi}(g_1^\alpha)$ and $u_2 = T_{\lambda, \Psi}(g_2^\alpha)$.

Let K be a fixed compact set of \mathbb{R}^d .

We clearly have, since $T_{\lambda, \Psi}$ is linear continuous on $L^1(K)$:

$$\|u_1\|_{L^1_x(K)} \leq C \|g_1^\alpha\|_1.$$

We notice that:

$$|\{(x, v), |g(x, v)| > \alpha\}| \leq \frac{1}{\alpha} \|g\|_{L^1} \leq \frac{1}{\alpha} C.$$

Since \mathcal{K} is equiintegrable, there exists $\alpha > 0$ such that for any $g \in \mathcal{K}$:

$$\int |g \mathbb{1}_{\{|g(x,v)| > \alpha\}}| dx dv \leq \frac{\eta}{C}.$$

Consequently for α large enough, we have:

$$\|u_1\|_{L_x^1(K)} \leq \eta.$$

The parameter α being fixed, we clearly see that $\{g_2^\alpha, g \in \mathcal{K}\}$ is a bounded subset of $L_{x,v}^1 \cap L_{x,v}^\infty$, and consequently of $L_{x,v}^2$. Because of Corollary 1, $\{u_2, u_2 = T_{\lambda,\Psi}(g_2^\alpha), g \in \mathcal{K}\}$ is relatively compact in $L_{loc}^2(dx)$. In particular it is relatively compact in $L_{loc}^1(dx)$.

As a result, we have shown that for any $\eta > 0$, there exists $\mathcal{K}_\eta \subset L_x^1(K)$ compact, such that $T_{\lambda,\Psi}(\mathcal{K}) \subset \mathcal{K}_\eta + B(0, \eta)$. So this family is precompact and consequently it is compact since $L_x^1(K)$ is a Banach space. \square

We deduce the preliminary result (which means that the first point implies the second in Theorem 1):

Theorem 3. *Let (f_ϵ) a family of $L_{loc}^1(dx \otimes dv)$ locally equiintegrable in x and v such that $(v \cdot \nabla_x f_\epsilon + F \cdot \nabla_v f_\epsilon)$ is a bounded family of $L_{loc}^1(dx \otimes dv)$. Then for all $\Psi \in \mathcal{C}_c^1(\mathbb{R}^d)$, the family $\rho_\epsilon(x) = \int f_\epsilon(x, v) \Psi(v) dv$ is relatively compact in $L_{loc}^1(dx)$.*

Proof. Let $\Psi \in \mathcal{C}_c^1(\mathbb{R}^d)$. Let $R > 0$ be a large number such that $\text{Supp } \Psi \subset B(0, R)$; we intend to show that $(\mathbb{1}_{B(0,R)}(x) \rho_\epsilon(x))$ is compact in $L^1(B(0, R))$. First of all, we can assume that the f_ϵ are compactly supported in the same compact set $K \subset \mathbb{R}^d \times \mathbb{R}^d$, with $B(0, R) \times B(0, R) \subset K$. Indeed we can multiply the family by a smooth function χ such that:

$$\begin{aligned} \text{supp } \chi &\subset K, \\ \chi &\equiv 1 \text{ on } B(0, R) \times B(0, R). \end{aligned}$$

We observe that :

$$\begin{aligned} v \cdot \nabla_x (\chi f_\epsilon) &= \chi (v \cdot \nabla_x f_\epsilon) + f_\epsilon (v \cdot \nabla_x \chi), \\ F \cdot \nabla_v (\chi f_\epsilon) &= \chi F \cdot \nabla_v (f_\epsilon) + f_\epsilon (F \cdot \nabla_v \chi). \end{aligned}$$

Thus the family (χf_ϵ) satisfies the same L^1 boundedness properties as (f_ϵ) . The equiintegrability property is also clearly preserved. Furthermore, for any x in $B(0, R)$, we have :

$$\int f_\epsilon(x, v) \Psi(v) dv = \int f_\epsilon(x, v) \chi(v) \Psi(v) dv$$

Consequently we are now in the case of functions supported in the same compact set.

We have for all $\epsilon > 0, \lambda > 0$, by linearity of the resolvent R_λ defined in Proposition 2 :

$$\begin{aligned} \int f_\epsilon(x, v) \Psi(v) dv &= \int R_\lambda (\lambda f_\epsilon + v \cdot \nabla_x f_\epsilon + F \cdot \nabla_v f_\epsilon) \Psi(v) dv \\ &= \lambda \int (R_\lambda f_\epsilon)(x, v) \Psi(v) dv + \int (R_\lambda (v \cdot \nabla_x f_\epsilon + F \cdot \nabla_v f_\epsilon))(x, v) \Psi(v) dv. \end{aligned}$$

$$\text{Let } \eta > 0. \text{ We take } \lambda = \sup_{\epsilon} \frac{\|(v \cdot \nabla_x f_\epsilon + F \cdot \nabla_v f_\epsilon)\|_{L_{x,v}^1} \|\Psi\|_{L^\infty}}{\eta}.$$

Then we have by Proposition 2 :

$$\begin{aligned} \left\| \int (R_\lambda(v \cdot \nabla_x f_\epsilon + F \cdot \nabla_v f_\epsilon))(x, v) \Psi(v) dv \right\|_{L_x^1} &\leq \|R_\lambda(v \cdot \nabla_x f_\epsilon + F \cdot \nabla_v f_\epsilon)\|_{L_{x,v}^1} \|\Psi\|_{L_v^\infty} \\ &\leq \frac{1}{\lambda} \|(v \cdot \nabla_x f_\epsilon + F \cdot \nabla_v f_\epsilon)\|_{L_{x,v}^1} \|\Psi\|_{L_v^\infty} \\ &\leq \eta. \end{aligned}$$

Moreover, since (f_ϵ) is bounded in $L^1(dx \otimes dv)$ and equiintegrable in x and v , Proposition 3 implies that the family $(\int R_\lambda(f_\epsilon) \Psi(v) dv)$ is relatively compact in $L_x^1(B(0, R))$. Finally we can argue as for the end of the proof of Proposition 3: for all $\eta > 0$, there exists $K_\eta \subset L_x^1(B(0, R))$ compact, such that $(\rho_\epsilon) \subset K_\eta + B(0, \eta)$. So this family is precompact and consequently it is compact since $L^1(B(0, R))$ is a Banach space. \square

2 Mixing properties of the operator $v \cdot \nabla_x + F \cdot \nabla_v$

2.1 Free transport case

In the case when $F = 0$, Bardos and Degond in [2] proved a mixing result (also referred to as a dispersion result for large time asymptotics) which is a key argument in the proof of Theorem 1 (with $F = 0$) by Golse and Saint-Raymond [10]. This kind of estimate was introduced for the study of classical solutions of the Vlasov-Poisson equation in three dimensions and for small initial data.

Lemma 2. *Let f be the solution to:*

$$\begin{cases} \partial_t f + v \cdot \nabla_x f = 0, \\ f(0, \cdot, \cdot) = f^0. \end{cases} \quad (2.1)$$

Then for all $t > 0$:

$$\|f(t)\|_{L_x^\infty(L_v^1)} \leq \frac{1}{|t|^d} \|f^0\|_{L_x^1(L_v^\infty)}. \quad (2.2)$$

For further results and related questions (Strichartz estimates...), we refer to Castella and Perthame [4] and Salort [15], [16], [17].

When f^0 is the indicator function of a set with "small" measure with respect to x , then the previous estimate (2.2) asserts that for $t > 0$, $f(t)$ is for any fixed x the indicator function of a set with a "small" measure in v (at least that we may estimate): this property is crucial for the following. In (2.2), there is blow-up when $t \rightarrow 0$, which is intuitive, but this does not matter since we have nevertheless a control of the left-hand side for any positive time.

Actually, for our purpose, parameter t is an artificial time (it does not have the usual physical meaning). It appears as an interpolation parameter in Lemma 4, and can be taken rather small. This is the reason why local in time mixing is sufficient. We will consequently look for local in time mixing properties. Anyway, the explicit study in Example 2 below shows that the dispersion inequality is in general false for large times when $F \neq 0$.

Proof of Lemma 2. The proof of this result is based on the explicit solution to (2.1), which is:

$$f(t, x, v) = f^0(x - tv, v)$$

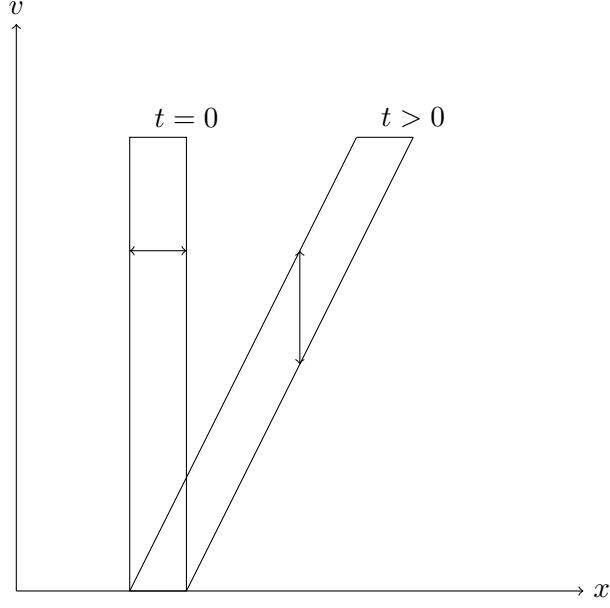


Figure 1: Mixing property for free transport

We now evaluate:

$$\begin{aligned}
\|f(t)\|_{L_x^\infty(L_v^1)} &= \sup_x \int f^0(x - tv, v) dv \\
&= \sup_x \int f^0(z, \frac{x - z}{t}) |t|^{-d} dz \\
&\leq |t|^{-d} \int \|f^0(z, \cdot)\|_\infty dz \\
&\leq |t|^{-d} \|f^0\|_{L_x^1(L_v^\infty)}.
\end{aligned}$$

The key argument is the change of variables $x - tv \mapsto z$, the jacobian of which is equal to t^{-d} . \square

We intend to do the same in the more complicated case when f is the solution of a transport equation with $F \neq 0$. Let us mention that in [2], Bardos and Degond actually prove the dispersion result for non zero force fields but with a polynomial decay in time. Here, this is not the case (the field F does not even depend on time t), but we will prove that the result holds anyway for small times.

2.2 Study of two examples

In the following examples, f is the explicit solution to the transport equation (1.4) with an initial condition f^0 and a force deriving from a potential.

Example 1

Force $F = -\nabla_x V$, with $V = -|x|^2/2$.

Let f be the solution to:

$$\begin{cases} \partial_t f + v \cdot \nabla_x f + x \cdot \nabla_v f = 0, \\ f(0, \cdot, \cdot) = f^0. \end{cases} \quad (2.3)$$

The effect of such a potential will be to make the particles escape faster to infinity. So we expect to have results very similar to those of lemma 2.

After straightforward computations we get :

$$f(t, x, v) = f^0 \left(x \left(\frac{e^t + e^{-t}}{2} \right) + v \left(\frac{e^{-t} - e^t}{2} \right), x \left(\frac{e^{-t} - e^t}{2} \right) + v \left(\frac{e^t + e^{-t}}{2} \right) \right),$$

which allows to show the same dispersion estimate with a factor $\frac{e^t - e^{-t}}{2}$ instead of t . For all $t > 0$, we have:

$$\|f(t)\|_{L_x^\infty(L_v^1)} \leq \frac{2^d}{(e^t - e^{-t})^d} \|f^0\|_{L_x^1(L_v^\infty)}. \quad (2.4)$$

Example 2

(Harmonic potential) Force $F = -\nabla_x V$, with $V = |x|^2/2$.

Let f be the solution to:

$$\begin{cases} \partial_t f + v \cdot \nabla_x f - x \cdot \nabla_v f = 0, \\ f(0, \cdot, \cdot) = f^0. \end{cases} \quad (2.5)$$

With such a potential, particles are expected to be confined and consequently do not drift to infinity. For this reason, it is hopeless to prove the analogue of Lemma 2 for large times (here there is no dispersion). As mentioned before, it does not matter since we only look for a result valid for small times. We expect that there is enough mixing in the phase space to prove the result.

After straightforward computations we explicitly have :

$$f(t, x, v) = f^0(x \cos t - v \sin t, v \cos t + x \sin t).$$

We observe here that the solution f is periodic with respect to time. Thus, as expected, it is not possible to prove any decay when $t \rightarrow +\infty$; nevertheless we can prove a mixing estimate with a factor $|\sin t|$ instead of t .

For all $t > 0$:

$$\|f(t)\|_{L_x^\infty(L_v^1)} \leq \frac{1}{|\sin t|^d} \|f^0\|_{L_x^1(L_v^\infty)}. \quad (2.6)$$

Of course, this estimate is useless when $t = k\pi, k \in \mathbb{N}^*$.

Remark 2. We notice that $\frac{e^t - e^{-t}}{2} \sim_0 t$ and $\sin(t) \sim_0 t$, which seems encouraging.

2.3 General case : F with Lipschitz regularity

The study of these two examples suggests that at least for small times, the mixing estimate is still satisfied, maybe with a corrector term which does not really matter.

One nice heuristic way to understand this is to see that since F is quite smooth, the dynamics associated to the operator $v \cdot \nabla_x + F \cdot \nabla_v$ is expected to be close to those of free transport, at least for small times.

Let $X(t; x, v)$ and $V(t; x, v)$ be the diffeomorphisms introduced in the method of characteristics in Section 1 and defined in (1.3).

Using Taylor's formula, we get by definition of X and V :

$$X(-t; x, v) = x - tv + \int_0^t (t-s) F(X(-s; x, v)) ds.$$

We recall Rademacher's theorem which asserts that $W^{1,\infty}$ functions are almost everywhere derivable. Hence, using Lebesgue domination theorem, we get:

$$\partial_v[X(-t; x, v)] = -tId + \int_0^t (t-s) \nabla_x F(X(-s; x, v)) \partial_v[X(-s; x, v)] ds. \quad (2.7)$$

We deduce the estimate :

$$\|\partial_v[X(-t; x, v)]\|_\infty \leq t + \int_0^t (t-s) \|\nabla_x F\|_\infty \|\partial_v[X(-s; x, v)]\|_\infty ds.$$

Gronwall's lemma implies then that:

$$\|\partial_v[X(-t; x, v)]\|_\infty \leq t e^{\frac{t^2}{2} \|\nabla_x F\|_\infty}. \quad (2.8)$$

We can also take the determinant of identity (2.7) :

$$\det(\partial_v[X(-t; x, v)]) = (-t)^d \det \left(Id - \frac{1}{t} \int_0^t (t-s) \nabla_x F(X(-s; x, v)) \partial_v[X(-s; x, v)] ds \right). \quad (2.9)$$

The right-hand side is the determinant of a matrix of the form $Id + A(t)$ where A is a matrix whose L^∞ norm is small for small times t (one can use estimate (2.8) to ensure that $\|A(t)\|_\infty = o(t)$). Consequently, in a neighborhood of 0, for any fixed x , $\partial_v[X(-t; x, v)]$ is invertible. Furthermore the map $v \mapsto X(-t; x, v)$ is injective for small positive times. Indeed, let $v \neq v'$. We compare:

$$X(-t; x, v') - X(-t; x, v) = t(v - v') + \int_0^t (t-s) [F(X(-s; x, v')) - F(X(-s; x, v))] ds.$$

Consequently we have:

$$|X(-t; x, v') - X(-t; x, v)| \leq t|v - v'| + \int_0^t (t-s) \|\nabla_x F\|_{L^\infty} |X(-s; x, v') - X(-s; x, v)| ds.$$

Thus, by Gronwall inequality we obtain:

$$|X(-t; x, v') - X(-t; x, v)| \leq t|v - v'| e^{\frac{t^2}{2} \|\nabla_x F\|_{L^\infty}}.$$

Finally we observe that:

$$\begin{aligned} |X(-t; x, v') - X(-t; x, v)| &\geq t|v - v'| - \left| \int_0^t (t-s) [F(X(-s; x, v')) - F(X(-s; x, v))] ds \right| \\ &\geq t|v - v'| - \int_0^t (t-s) \|\nabla_x F\|_{L^\infty} |X(-s; x, v') - X(-s; x, v)| ds \\ &\geq |v - v'| \left(t - \int_0^t (t-s) s e^{\frac{s^2}{2} \|\nabla_x F\|_{L^\infty}} \|\nabla_x F\|_{L^\infty} ds \right). \end{aligned}$$

Consequently, there is a maximal time $\tau_0 > 0$, depending only on $\|\nabla_x F\|_{L^\infty}$ such that for any $|t| \leq \tau_0$, we have :

$$|X(-t; x, v') - X(-t; x, v)| \geq \frac{t}{2} |v - v'|.$$

This proves our claim.

Thus, by the local inversion theorem, this map is a \mathcal{C}^1 diffeomorphism on its image.

We have now the following elementary quantitative estimate :

Lemma 3. *Let $t \mapsto A(t)$ be a continuous map defined on a neighborhood of 0, such that $\|A(t)\|_\infty = o(t)$. Then for small times:*

$$\det(\text{Id} + A(t)) \geq 1 - d!\|A(t)\|_\infty.$$

We recall that d is the space dimension and $d! = 1 \times 2 \times \dots \times d$.

We apply this lemma to (2.9), which allows us to say that there exists a maximal time $\tau > 0$ such that for any $|t| \leq \tau$, we have :

$$|\det(\partial_v[X(-t; x, v)])|^{-1} \leq 2|t|^{-d}. \quad (2.10)$$

We have proved that $v \mapsto X(-t; x, v)$ is a \mathcal{C}^1 diffeomorphism such that the jacobian of its inverse satisfies (2.10) in a neighborhood of $t = 0$. We can consequently conclude as in the proof of Lemma 2 (by performing the change of variables $X(-t; x, v) \mapsto v$).

As a result we have proved the proposition :

Proposition 4. *Let $F(x)$ be a Lipschitz vector field. There exists a maximal time $\tau > 0$ (depending only on $\|\nabla_x F\|_{L^\infty}$) such that, if f is the solution to the transport equation:*

$$\begin{cases} \partial_t f + v \cdot \nabla_x f + F \cdot \nabla_v f = 0, \\ f(0, \cdot, \cdot) = f^0 \in L^p(dx \otimes dv). \end{cases} \quad (2.11)$$

Then:

$$\forall |t| \leq \tau, \|f(t)\|_{L_x^\infty(L_v^1)} \leq \frac{2}{|t|^d} \|f^0\|_{L_x^1(L_v^\infty)}. \quad (2.12)$$

Remark 3. *If one writes down more explicit estimates, it can be easily shown that τ is bounded from below by T defined as the only positive solution to the equation:*

$$\frac{d!}{3} \|\nabla_x F\|_\infty T^2 e^{\|\nabla_x F\|_\infty \frac{T^2}{2}} = 1. \quad (2.13)$$

Remark 4. *Of course, one can replace the factor 2 in the mixing estimate by any $q > 1$ (and the maximal time τ will depend also on q).*

3 From local equiintegrability in velocity to local equiintegrability in position and velocity

In this section, we finally proceed as in [10], with some slight modifications adapted to our case. We start from the following Green's formula :

Lemma 4. *Let $f \in L^1(dx \otimes dv)$ with compact support such that $v \cdot \nabla_x f + F \cdot \nabla_v f \in L^1(dx \otimes dv)$. Then for all $\Phi^0 \in L^\infty(dx \otimes dv)$, we have for all $t \in \mathbb{R}_+^*$:*

$$\begin{aligned} \int f(x, v) \Phi^0(x, v) dx dv &= \int f(x, v) \Phi(t, x, v) dx dv \\ &\quad - \int_0^t \int \Phi(s, x, v) (v \cdot \nabla_x f + F \cdot \nabla_v f) ds dx dv, \end{aligned} \quad (3.1)$$

where Φ is the solution to:

$$\begin{cases} \partial_t \Phi + v \cdot \nabla_x \Phi + F \cdot \nabla_v \Phi = 0 \\ \Phi|_{t=0} = \Phi_0. \end{cases} \quad (3.2)$$

Proof. We have for all $t > 0$, $\int_{\Omega} f(x, v) (\partial_t + v \cdot \nabla_x + F \cdot \nabla_v) \Phi(s, x, v) ds dx dv = 0$, where $\Omega =]0, t[\times \mathbb{R}^d \times \mathbb{R}^d$. We first have:

$$\int_{\Omega} f(x, v) \partial_t \Phi(s, x, v) ds dx dv = \int \int f(x, v) \Phi(t, x, v) dx dv - \int \int f(x, v) \Phi^0(x, v) dx dv.$$

Finally, by Green's formula we obtain:

$$\begin{aligned} \int_0^t \int f(x, v) (v \cdot \nabla_x + F \cdot \nabla_v) \Phi(s, x, v) ds dx dv \\ = - \int_0^t \int \Phi(s, x, v) (v \cdot \nabla_x + F \cdot \nabla_v) f(x, v) ds dx dv. \end{aligned}$$

There is no contribution from the boundaries since f is compactly supported. \square

Lemma 5. *Let (f_{ϵ}) a bounded family of $L^1_{loc}(dx \otimes dv)$ locally integrable in v such that $(v \cdot \nabla_x f_{\epsilon} + F \cdot \nabla_v f_{\epsilon})$ is a bounded family of $L^1_{loc}(dx \otimes dv)$. Then for all $\Psi \in \mathcal{C}_c^1(\mathbb{R}^d)$, such that $\Psi \geq 0$, the family $\rho_{\epsilon}(x) = \int |f_{\epsilon}(x, v)| \Psi(v) dv$ is locally equiintegrable.*

Proof. Let K_1 be a compact subset of \mathbb{R}^d . We want to prove that $(\mathbb{1}_{K_1} \rho_{\epsilon}(x))$ is equiintegrable. Without loss of generality, we can assume as previously that the f_{ϵ} are supported in the same compact support $K = K_1 \times K_2$ and that Ψ is compactly supported in K_2 . Furthermore, the formula $\nabla |f_{\epsilon}| = \text{sign}(f_{\epsilon}) \nabla f_{\epsilon}$ shows that the $|f_{\epsilon}|$ satisfy the same assumptions of equiintegrability and L^1 boundedness as the family (f_{ϵ}) . For the sake of readability, we will thus assume that f_{ϵ} are almost everywhere non-negative instead of considering $|f_{\epsilon}|$. Finally we may assume that $\|\Psi\|_{\infty} = 1$ (multiplying by a constant does not change the equiintegrability property).

The idea of the proof is to show that thanks to the mixing properties established previously, the equiintegrability in v provides some equiintegrability in x .

Let $\eta > 0$. By definition of the local equiintegrability in v , we obtain a parameter $\alpha > 0$ associated to K and η . We also consider parameters $\alpha' > 0$ and $t \in]0, \tau[$ (where τ is the maximal time in Proposition 4) to be fixed ultimately. We mention that t will be chosen only after α' is fixed.

Let A a bounded measurable subset included in K_1 with $|A| \leq \alpha'$. We consider $\Phi^0(x, v) = \mathbb{1}_A(x)$ and Φ the solution of the transport equation (3.2) with Φ^0 as initial data.

Observe now that we have $\|\Phi^0\|_{L^1_x(L^{\infty}_v)} = |A|$. Moreover, since Φ^0 takes its values in $\{0, 1\}$, it is also the case for Φ (this is a plain consequence of the transport of the data).

We define for all $s > 0$ and for all $x \in \mathbb{R}^d$, the set $A(s)_x = \{v \in \mathbb{R}^d, \Phi(s, x, v) = 1\}$. At this point of the proof, we make a crucial use of the mixing property stated in Proposition 4 :

$$\begin{aligned} \sup_x |A(t)_x| &= \sup_x \int \Phi(t, x, v) dv \\ &= \|\Phi(t, \cdot, \cdot)\|_{L^{\infty}_x(L^1_v)} \\ &\leq 2|t|^{-d} \underbrace{\|\Phi^0\|_{L^1_x(L^{\infty}_v)}}_{|A| \leq \alpha'} \\ &\leq \alpha, \end{aligned}$$

if we choose α' satisfying $\alpha' < \frac{1}{2}t^D\alpha$.

Thanks to Lemma 4:

$$\begin{aligned} \int f(x, v)\Psi(v)\Phi^0(x, v)dx dv &= \int f(x, v)\Psi(v)\Phi(t, x, v)dx dv \\ &\quad - \int_0^t \int \Phi(s, x, v)(v \cdot \nabla_x + F \cdot \nabla_v)(f_\epsilon(x, v)\Psi(v))dx dv ds. \end{aligned}$$

In other words, the operator $\partial_t + v \cdot \nabla_x + F \cdot \nabla_v$ has transported the indicator function and has transformed a subset small in x into a subset small in v .

By definition of ρ_ϵ , we have:

$$\int f(x, v)\Psi(v)\Phi^0(x, v)dx dv = \int \mathbb{1}_A(x)\rho_\epsilon(x)dx.$$

By definition of $A(t)_x$ we also have:

$$\int f(x, v)\Psi(v)\Phi(t, x, v)dx dv = \int \left(\int_{A(t)_x} f_\epsilon(x, v)\Psi(v)dv \right) dx.$$

Thus, since (f_ϵ) are locally equiintegrable in v we may evaluate:

$$\begin{aligned} \int \left(\int_{A(t)_x} f_\epsilon(x, v)\Psi(v)dv \right) dx &\leq \int \int_{A(t)_x} |f_\epsilon| \mathbb{1}_K \underbrace{\|\Psi\|_\infty}_{=1} dx dv \\ &\leq \eta. \end{aligned}$$

Finally we have:

$$\begin{aligned} \int \mathbb{1}_A(x)\rho_\epsilon(x)dx &= \int \left(\int_{A(t)_x} f_\epsilon(x, v)\Psi(v)dv \right) dx \\ &\quad - \int_0^t \int \Phi(s, x, v)(v \cdot \nabla_x + F \cdot \nabla_v)(f_\epsilon(x, v)\Psi(v))ds dx dv \\ &\leq \eta + \int_0^t \int |\Phi(s, x, v)| |(v \cdot \nabla_x + F \cdot \nabla_v)(f_\epsilon(x, v)\Psi(v))| ds dx dv \\ &\leq \eta + t \left[\underbrace{\|\Psi\Phi\|_\infty}_{\leq 1} \|v \cdot \nabla_x f_\epsilon + F \cdot \nabla_v f_\epsilon\|_1 + \underbrace{\|\Phi\|_\infty}_{=1} \|F \cdot \nabla_v \Psi(v)\|_\infty \|f_\epsilon\|_1 \right] \\ &\leq 2\eta, \end{aligned}$$

by taking t sufficiently small:

$$t < \frac{\eta}{\sup_\epsilon \|v \cdot \nabla_x f_\epsilon + F \cdot \nabla_v f_\epsilon\|_1 + \|F \cdot \nabla_v \Psi(v)\|_\infty \|f_\epsilon\|_1}.$$

This finally proves that (ρ_ϵ) is locally equiintegrable in x . \square

Lemma 6. *Let (g_ϵ) a bounded family of $L^1_{loc}(dx \otimes dv)$ locally equiintegrable in v . If for all $\Psi \in \mathcal{C}_c^1(\mathbb{R}^d)$ such that $\Psi \geq 0$, $x \mapsto \int |g_\epsilon(x, v)|\Psi(v)dv$ is locally equiintegrable (in x), then (g_ϵ) is locally equiintegrable in x and v .*

Proof. Let K be a compact subset of $\mathbb{R}^d \times \mathbb{R}^d$. We want to prove that $(\mathbb{1}_K g_\epsilon)$ is equiintegrable in x and v . As before, we can clearly assume that the g_ϵ are compactly supported in K .

Let $\eta > 0$. By definition of the local equiintegrability in v for (g_ϵ) , we obtain $\alpha_1 > 0$ associated to η and K .

Let $\Psi \in \mathcal{C}_c^1(\mathbb{R}^d)$ a smooth non-negative and compactly supported function such that $\Psi \equiv 1$ on $p_v(K)$ (where $p_v(K)$ is the projection of K on \mathbb{R}_v^d). By assumption, there exists $\alpha_2 > 0$ such that for any $A \subset \mathbb{R}^d$ measurable set satisfying $|A| \leq \alpha_2$,

$$\int_A \left(\int |g_\epsilon| \Psi dv \right) dx < \eta.$$

Let B a measurable subset of $\mathbb{R}^d \times \mathbb{R}^d$ such that $|B| < \inf(\alpha_1^2, \alpha_2^2)$. We define for all $x \in \mathbb{R}^d$, $B_x = \{v \in \mathbb{R}^d, (x, v) \in B\}$.

We consider now $E = \{x \in \mathbb{R}^d, |B_x| \leq |B|^{1/2}\}$: this is the subset of x for which there exist few v such that $(x, v) \in B$. Consequently for this subset, we can use the local equiintegrability in v .

Concerning $B \setminus E$, on the contrary, we can not use this property, but thanks to Chebychev's inequality we show that this subset is of small measure, which allows us to use this time the local equiintegrability in x of $\int |g_\epsilon(x, v)| \Psi(v) dv$:

$$\begin{aligned} |E^c| &= |\{x \in \mathbb{R}^d, |B_x| > |B|^{1/2}\}| \\ &\leq \frac{|B|}{|B|^{1/2}} \\ &\leq \alpha_2. \end{aligned}$$

Hence we have :

$$\begin{aligned} \int \mathbb{1}_B |g_\epsilon| dx dv &\leq \int_E \left(\int_{B_x} |g_\epsilon| dv \right) dx + \int_{E^c} \left(\int |g_\epsilon| dv \right) dx \\ &\leq \eta + \int_{E^c} \left(\int |g_\epsilon| \Psi(v) dv \right) dx \\ &\leq 2\eta. \end{aligned}$$

This shows the expected result. □

We are now able to conclude the proof of Theorem 1.

End of the proof of Theorem 1. If we successively apply Lemmas 5 and 6, we deduce that the family (f_ϵ) is locally equiintegrable in x and v .

Finally we have shown in Section 1 that the first point implies the second. □

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